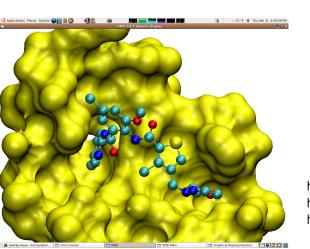


## Chemical Physics II



#### More Stat. Thermo Kinetics Protein Folding

. . .



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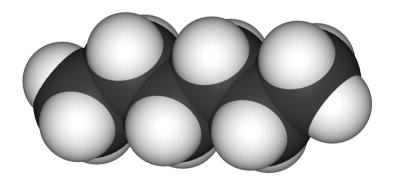
# Continuing Theme

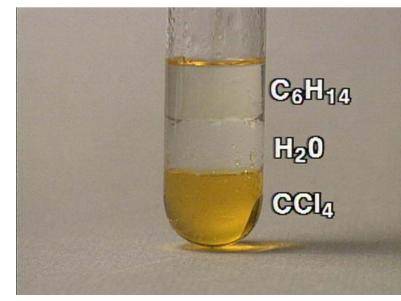
How can we understand the bulk properties of a substance from the properties of individual molecules

H<sub>3</sub>C-CH<sub>3</sub>-CH<sub>3</sub>-CH<sub>3</sub>-CH<sub>3</sub>-CH<sub>3</sub>

?

Density=0.6548 g/ml Boiling point=69 °C Melting point=-95°C etc.



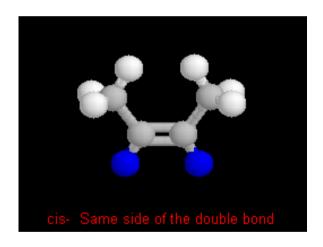


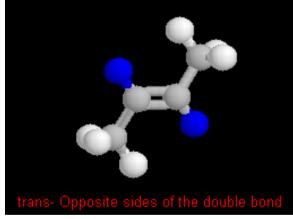
#### The Game Plan

- Review the two-state system
  - butene cis/trans example
  - activation energy and kinetics
- Brief intermezzo on the free energies
- Generalize the two-state problem to an N-state problem
  - We will see how this relates to the folding of proteins
  - This will give a new perspective on entropy

#### **Butene Gas**

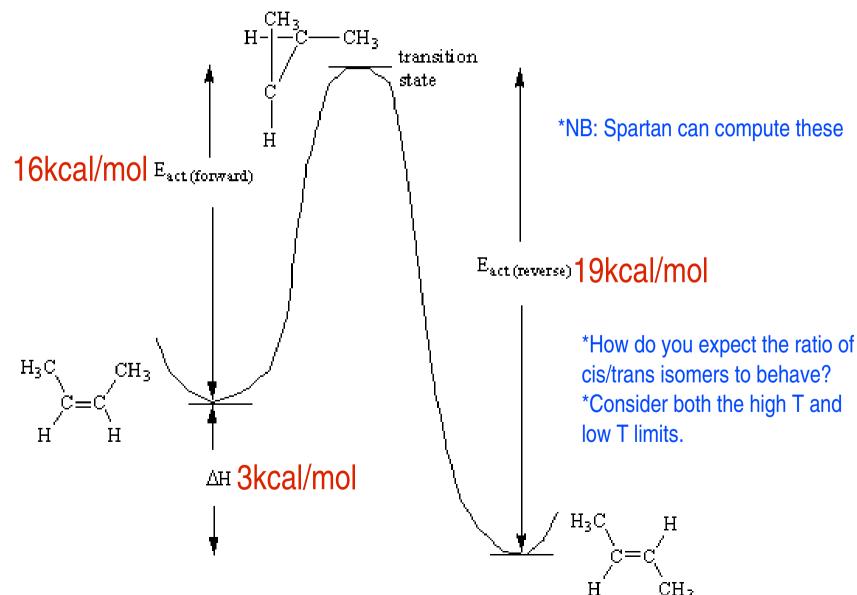
 Recall, some molecules of the gas will be in the cis conformation, others will be in the trans conformation:





#### **Butene Gas**

The butene potential energy surface



http://usm.maine.edu/~newton/CHY100/Transparencies/Chapter7/CoordinateDiagram.gif

#### Butene Gas cis/trans Probability

Three key properties of our probability function:

1. It should be a function of energy and temperature

$$P_c = P_c(E_c, T)$$
 and  $P_t = P_t(E_t, T)$ 

2. Independent events should multiply, ie for any two butene molecules

$$P_{c1,c2} = P_c \cdot P_c$$
;  $P_{c1,t2} = P_{t1,c2} = P_c \cdot P_t$ ;  $P_{t1,t2} = P_t \cdot P_t$ 

3. Probabilities should be normalized to 1

$$P_c + P_t = 1$$

Three key properties of our probability function:

1. It should be a function of energy and temperature

$$P_c \propto e^{-E_c/RT}$$
 and  $P_t \propto e^{-E_t/RT}$ 

$$dent events should multiply$$
Boltzmann
Factor

2. Independent events should multiply

$$P_{c1,t2} = P_c \cdot P_t \propto e^{-E_c/RT} \cdot e^{-E_t/RT} = e^{-(E_c + E_t)/RT}$$

Probabilities should be normalized to 1

$$\begin{aligned} P_c + P_t = 1 = & \frac{e^{-E_c/RT}}{Q} + \frac{e^{-E_t/RT}}{Q} = & \frac{e^{-E_c/RT} + e^{-E_t/RT}}{Q} \\ \text{where } Q = & e^{-E_c/RT} + e^{-E_t/RT} \end{aligned}$$

#### Butene Gas cis/trans Probability

So our final result was,

$$P_{c} = \frac{e^{-E_{c}/RT}}{Q} = \frac{e^{-E_{c}/RT}}{e^{-E_{c}/RT} + e^{-E_{t}/RT}}$$

$$P_{t} = \frac{e^{-E_{t}/RT}}{Q} = \frac{e^{-E_{t}/RT}}{e^{-E_{c}/RT} + e^{-E_{t}/RT}}$$

#### Butene Gas Isomerization Kinetics

- We estimated the rate at which the a jug of all cis butene will isomerize to trans butene by finding
  - what fraction of the molecules at a given time are at the transition state
  - this is the same calculation we just did, but now ∆E=16kcal/mol instead of 3kcal/mol

$$k_{cis \to trans} \propto P(tst) s^{-1} = \frac{e^{\frac{-\Delta E_{ctst}}{RT}}}{\frac{-\Delta E_{ctst}}{1 + e^{\frac{-\Delta E_{ctst}}{RT}}}} s^{-1} = \frac{e^{\frac{-16kcal/mol}{(.002 \, kcal/mol \cdot K)300K}}}{\frac{-16kcal/mol}{1 + e^{\frac{-16kcal/mol}{(.002 \, kcal/mol \cdot K)300K}}}} s^{-1}$$

$$\Rightarrow k_{cis \to trans} \propto 10^{-12}/s$$
 ie 30000yr halflife

## Free Energy Intermezzo

From last time,

$$\Delta G = \Delta H - T \Delta S$$

$$\Delta H = \Delta E + P \Delta V$$
, but  $P \Delta V \approx 0$  for bio.  
 $\Rightarrow \Delta H \approx \Delta E$ 

$$\Rightarrow \Delta G \approx \Delta E - T \Delta S \equiv \Delta A$$

Gibbs Free Energy

Helmholtz Free Energy (More important too us than it might look at first glance.)

## Free Energy Intermezzo

- We have now defined,
  - The partition function: Q
  - The Helmholtz free energy: ΔA
- If we can relate  $\Delta A$  and Q, then
  - can express the free energy in terms of Boltzmann factors
  - can derive other thermodynamic parameters such as pressure, entropy, heat capacity, isothermal compressibility, etc

#### Free Energy Intermezzo

It can be shown,

$$\Delta A = -RT \ln(Q)$$

$$\Pr = RT \frac{\partial \ln Q}{\partial V}$$

$$\langle E \rangle = RT^2 \frac{\partial \ln Q}{\partial T}$$

$$S = \frac{\partial}{\partial T} (RT \ln Q)$$

etc...

#### Generalizing the 2-State Solution

- We have now seen that we can think of our butene molecule rattling between three different states
  - The trans state (lowest in energy, most stable)
  - The cis state (higher in energy, but still long lived)
  - The transition state (very high in energy, short lived)
- We should be able to derive the the probability of a butene molecule occupying any of these three states

#### Generalizing the 2-State Solution

For 3 states we find,

$$P_{t} = \frac{e^{-E_{t}/RT}}{Q} = \frac{e^{-E_{t}/RT}}{e^{-E_{c}/RT} + e^{-E_{t}/RT} + e^{-E_{tst}/RT}} = .99_{298K} ; .82_{1000K}$$

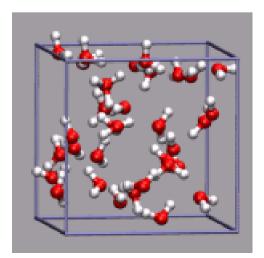
$$P_{c} = \frac{e^{-E_{c}/RT}}{Q} = \frac{e^{-E_{c}/RT}}{e^{-E_{c}/RT} + e^{-E_{t}/RT} + e^{-E_{tst}/RT}} = .01_{298K} ; .18_{1000K}$$

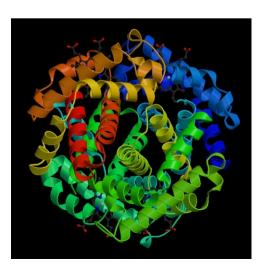
$$P_{tst} = \frac{e^{-E_{tst}/RT}}{Q} = \frac{e^{-E_{tst}/RT}}{e^{-E_{c}/RT} + e^{-E_{t}/RT} + e^{-E_{tst}/RT}} = 10^{-12}_{298K} ; 10^{-4}_{1000K}$$

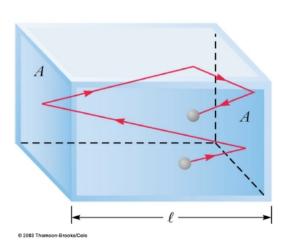
where  $E_t = 0 \text{kcal/mol}$ ;  $E_c = 3 \text{kcal/mol}$ ;  $E_{tst} = 16 \text{kcal/mol}$ 

#### Many State Problems

- Most "real world" chemical problems have very many states:
  - The number of ways liquid water molecules can arrange themselves in a cup
  - The number different shapes a protein can trace
  - The number of different positions and velocities gas particles can assume in a box



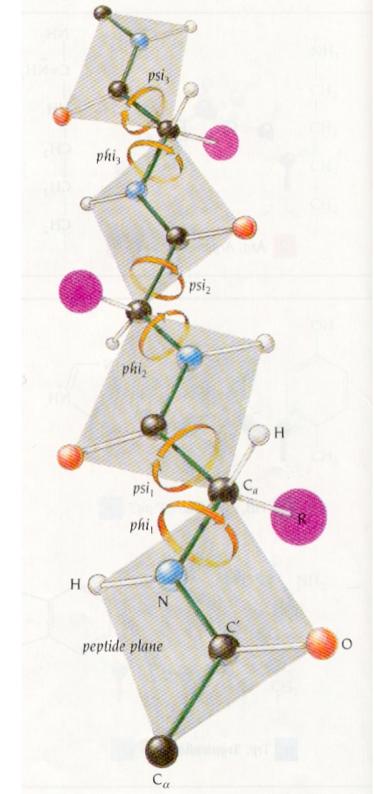




http://www.theochem.ruhr-uni-bochum.de/~axel.kohlmeyer/cpmd-vmd/unitcell-movie.gif

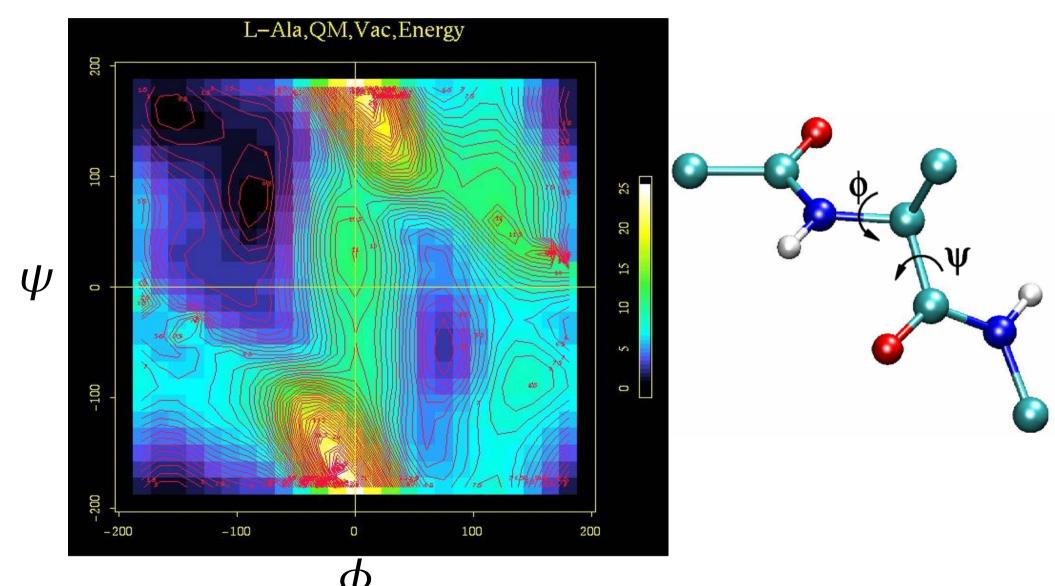
http://iwrwww1.fzk.de/biostruct/Assets/1a00x500.jpg

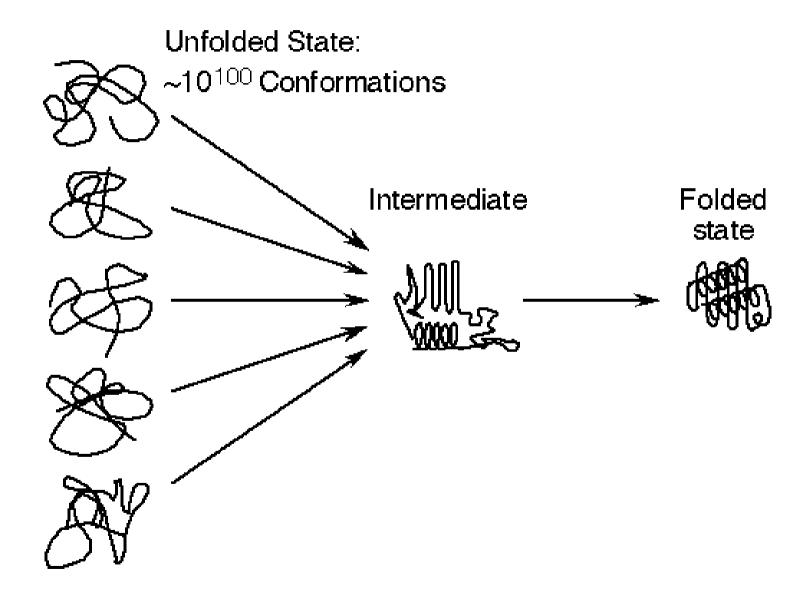
- Proteins are polymers
  - think of this as a string of beads
  - each bead is an amino acid
  - each amino acid has two rotatable bonds (phi, psi) connecting it to each other amino acid
  - the string changes shape by rotating around these bonds
  - When in water, each protein will spontaneously fold into a specific, well defined, shape



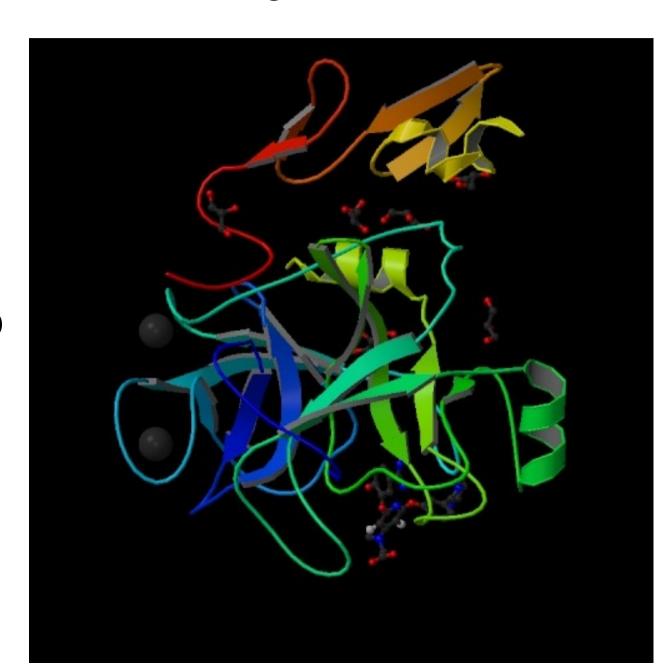
- Most proteins contain 100-200 amino acids
- Each amino acid can assume roughly ten different pairs of phi and psi angles
- So the total conformational space (ie all the shapes it can fold into) is 10<sup>100</sup> states
- Our goal is to understand the thermodynamics of protein folding

Visualization of the ~10 stable states for each amino acid





- The single folded state
   of a given protein is
   typically quite
   complicated
- It must be in this fold to function
- Prion diseases are caused by unfolded proteins



- This is a many state problem
  - There exists 1 single folded state with energy E<sub>F</sub>
  - There exist ~10<sup>100</sup> unfolded states with energy E<sub>...</sub>
- We will treat each of the unfolded states as though they have the same energy
- We will see
  - the large number of unfolded states opposes folding (entropy)
  - the favorable interaction energy of the folded state promotes folding (enthalpy)

We find for protein being folded or unfolded,

$$Q = e^{-E_F/RT} + 10^{100} e^{-E_U/RT} = e^{-E_F/RT} + \Omega e^{-E_U/RT}$$

$$P_F = \frac{e^{-E_F/RT}}{Q} = \frac{e^{-E_F/RT}}{e^{-E_F/RT} + 10^{100} e^{-E_U/RT}}$$
"Degeneracy"=10<sup>100</sup>

$$P_{U} = 10^{100} \frac{e^{-E_{U}/RT}}{Q} = \frac{e^{-E_{U}/RT}}{e^{-E_{F}/RT} + 10^{100} e^{-E_{U}/RT}}$$

What can we learn about protein folding from this solution?

- Experimentally,
  - Most proteins are nearly entirely folded at 25°C
  - Most proteins are half unfolded at 75°C
  - Most proteins are nearly entirely unfolded at 90°C
- Notice that this process is much more temperature dependent than the butene isomerization
- For the rest of the discussion we will choose our energy scale such that  $E_{E}=0$  and  $E_{II}=\Delta E_{II}$

We can calculate  $\Delta E_U$  by forcing  $P_F = .5$  @ 350K,

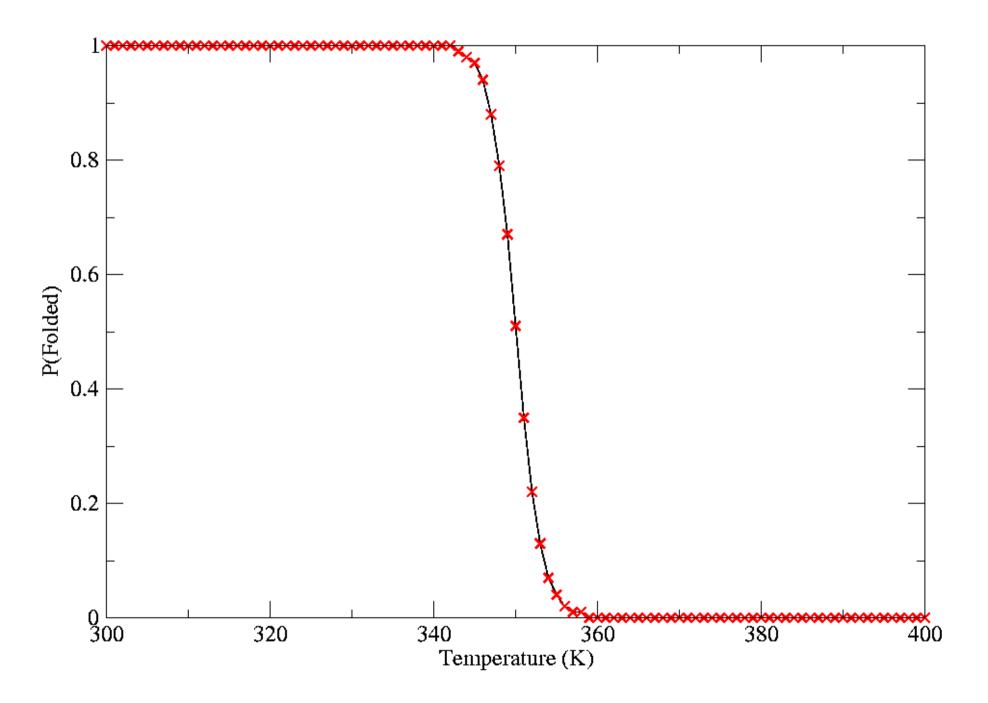
$$P_{F} = \frac{e^{-E_{F}/RT}}{e^{-E_{F}/RT} + 10^{100} e^{-E_{U}/RT}} = \frac{1}{1 + 10^{100} e^{-\Delta E_{U}/R \cdot 350K}} \equiv .5$$

$$\Rightarrow \Delta E_U = 161.2 \text{kcal/mol}$$

What does this say about the occupancy of the folded state at different temperatures?

$$P_{F@300K} = \frac{1}{1 + 10^{100} e^{\frac{-161.2 \text{kcal/mol}}{R \cdot 300K}}} \approx 1 \text{ and } P_{F@390K} \approx 10^{-11}$$

Probability a generic protein is folded as a function of temperature



We can calculate  $\Delta$  A of unfolding from this derived data,

$$\Delta A_{F \to U} = -RT \ln \frac{Q_U}{Q_F} = -RT \ln \left[ \frac{\Omega \cdot e^{-\Delta E_U/RT}}{1} \right]$$

$$\Rightarrow \Delta A_{F \to U} = -RT \ln e^{-\Delta E_U/RT} - RT \ln \Omega$$

$$\Rightarrow \Delta A_{F \to U} = \Delta E_U - RT \ln \Omega$$

but we have already showed,

$$\Delta A_{F \to U} = \Delta E_U - T \cdot \Delta S \Rightarrow \Delta S_U = R \ln \Omega$$

The value of  $\Delta$  A at 300K is also interesting,

$$\Delta A_{F \to U} = -RT \ln \left[ \Omega \cdot e^{-\Delta E_U/RT} \right]$$

$$\Rightarrow \Delta A_{F \to U} = \Delta E_U - RT \ln \Omega$$

$$\Rightarrow \Delta A_{F \to U} = 161.2 \text{kcal/mol} - 138.2 \text{kcal/mol} = 23 \text{kcal/mol}$$

For perspective, a water – water hydrogen bond is 5kcal/mol

#### Conclusions

- Proteins have a relatively small and abrupt window of thermal stability
  - this is because both the entropic and enthalpic terms are very large and opposed
- These methods of calculating  $\Delta E$ ,  $\Delta A$ , and  $\Delta S$  are at the core of computational chemistry
- We typically get ∆E from a model
- Once we have ΔE for each state we can calculate bulk thermodynamic properties