

Using AdS/CFT to visualize jet and wake-substructure in heavy ion collisions



[Abstract and more info](#)

Arjun Kudinoor

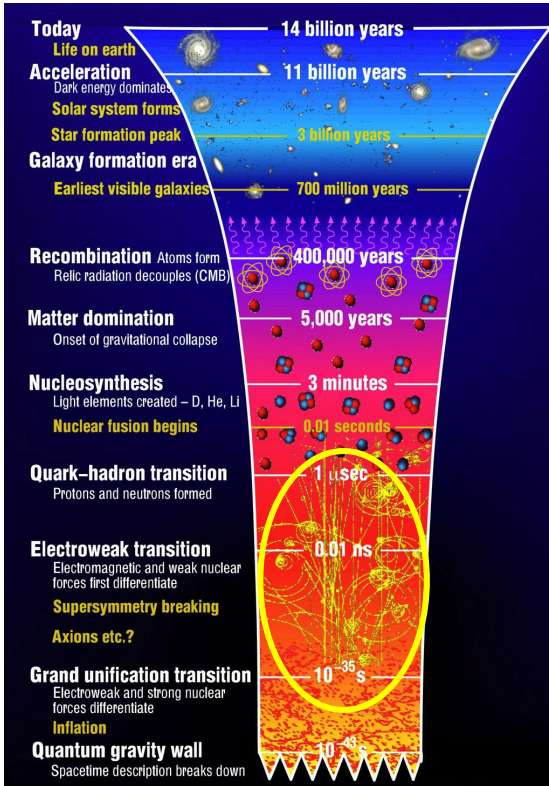
Collaborators: Krishna Rajagopal (MIT), Daniel Pablos (INFN)

- **Quark gluon plasma and jets**
- **Using AdS/CFT as a useful nonperturbative technique**
- **Visualizing jet and wake substructure**
- **Results from gamma-jet calculations**

TOPIC 1

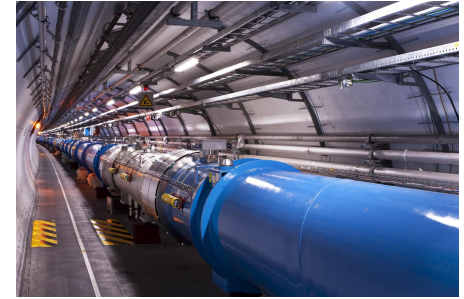
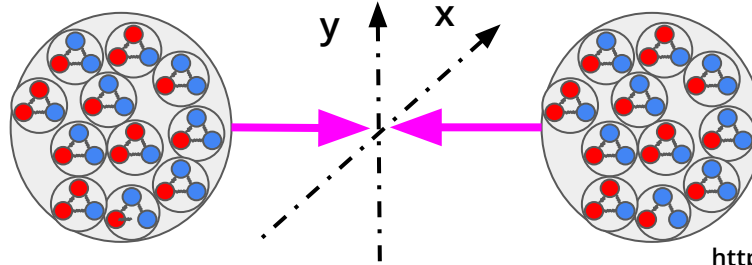
Quark gluon plasma and jets

1. THE EARLY UNIVERSE



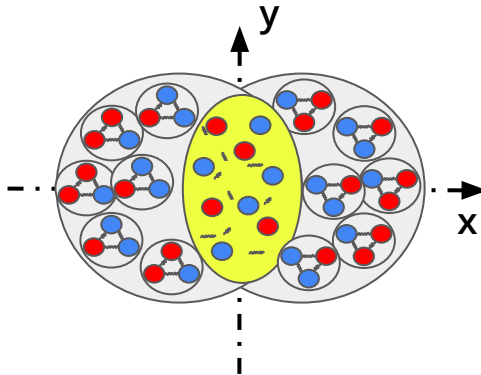
From: Centre for Theoretical
Cosmology, Cambridge University

2. HEAVY ION COLLISIONS



https://en.wikipedia.org/wiki/Large_Hadron_Collider

3. QUARK GLUON PLASMA

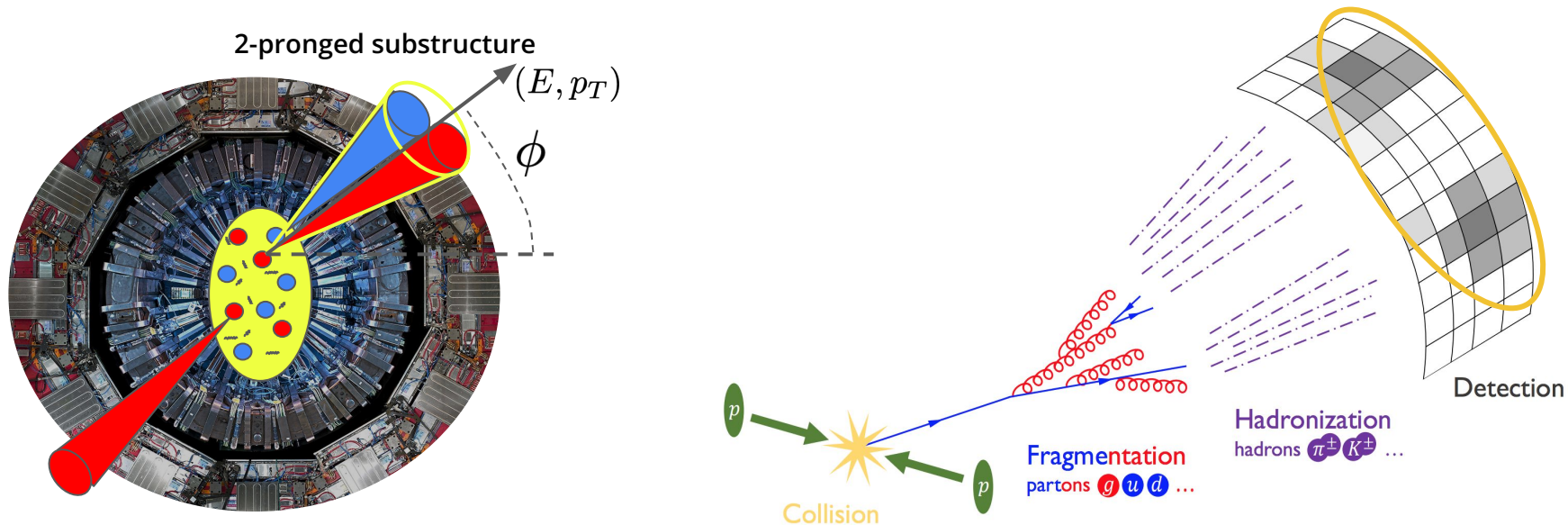


- The **first liquid** to exist
- Origin of protons and neutrons
- At a few trillion degrees, the **hottest liquid** that has ever existed
- The **most liquid liquid** to exist
- Quarks are **deconfined**, yet **strongly interacting**

https://indico.fysik.su.se/event/8016/contributions/12121/attachments/5069/6726/QC23_Rajagopal_Lectures_v2.pdf

Often, energetic sprays of hadrons shoot out of the QGP. We call these sprays **jets**. Jets are probes to understand the physics of QGP because they carry information about **energy, momentum, angular distribution, and internal structure** that can reveal the nature of QCD physics in quark gluon plasma.

Example: **Jet energy loss** can tell us about the **viscosity** and **density** of QGP

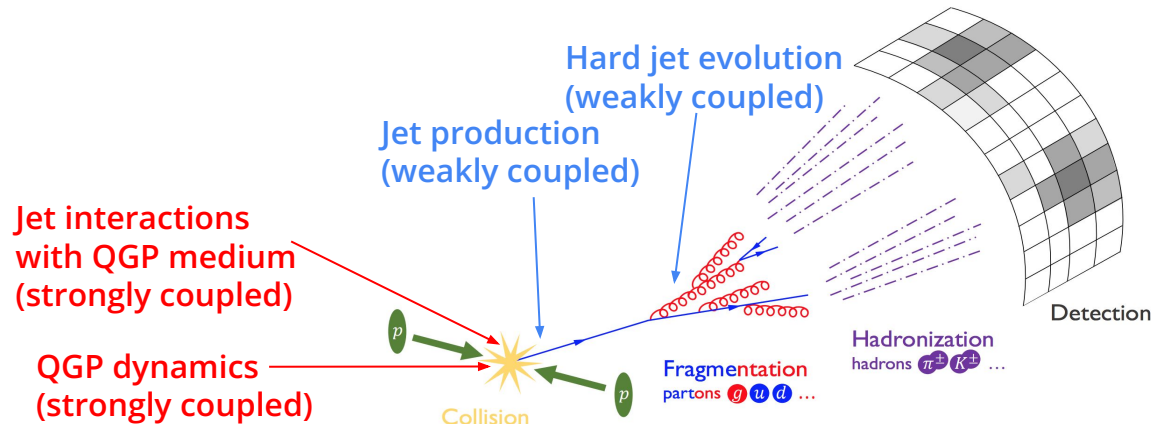


Detector image from: CMS

Seems simple enough... What's the catch?

The physics of jets and QGP hydrodynamics have both **weakly** and **strongly** coupled aspects. Calculations are **intractable** at strong coupling using standard perturbative methods.

“A successful phenomenological model that describes the modifications of jets in the medium, today, must be a **hybrid model** in which one can simultaneously treat the weakly coupled physics of **jet production** and **hard jet evolution** and the strongly coupled dynamics of **the [QGP] medium** and the **soft exchanges between the jet and the medium**” (arXiv:1405.3864v3)



**How do we perform jet energy loss
calculations at strong coupling?**

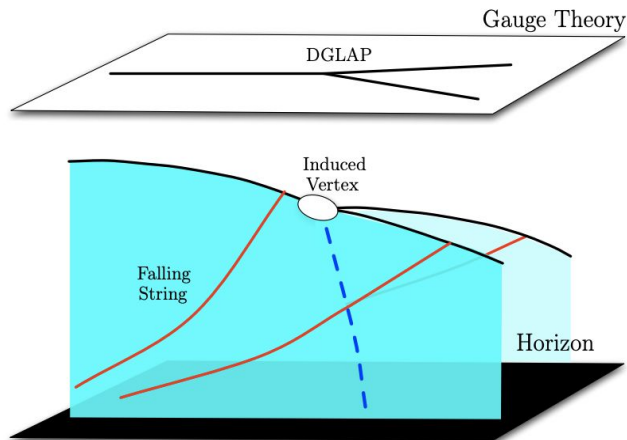
TOPIC 2

Using AdS/CFT

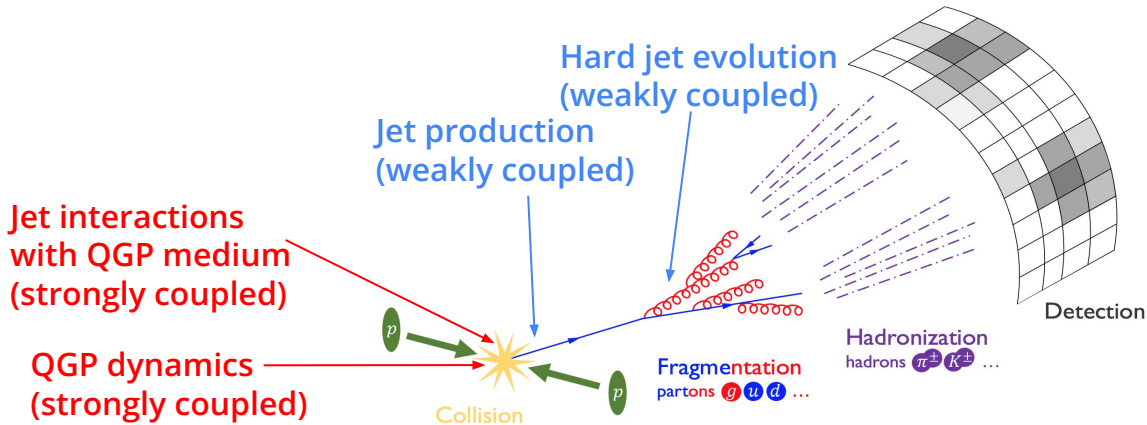
A useful nonperturbative technique

- Treat weakly coupled physics perturbatively
- Treat strongly coupled processes using AdS/CFT
 - Find the stringy gravity dual of QCD ← **umm...**
 - Describe your particles in using strings that hang from the boundary theory into the bulk spacetime
 - Calculate the observables you desire (energy loss, momenta, etc.)

Difficult calculations in **strongly coupled gauge theories** may be solved in its more tractable **weakly coupled (small curvature) gravitational dual**.



arXiv:1405.3864v3



<https://www.ericmetodiev.com/post/jetformation/>

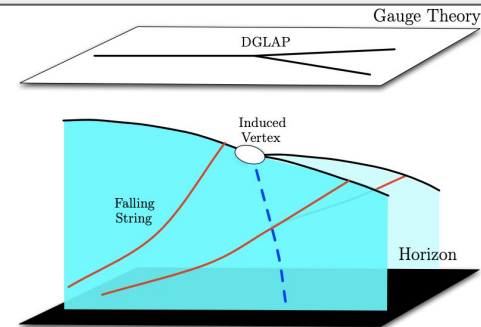
Use an $\mathcal{N} = 4$ $SU(N_c)$ SYM theory instead! The hot strongly coupled liquid phases of $\mathcal{N} = 4$ $SU(N_c)$ SYM theory and QCD are more similar to each other than their vacua and low energy physics (the problematic energy sector that contributes to QCD's nonconformality).

Differences between QCD and $\mathcal{N} = 4$ SYM include

- $N_c = 3$ for QCD, whereas we take the $N_c \rightarrow \infty$ limit for $\mathcal{N} = 4$ SYM calculations
- QCD is not conformal, whereas $\mathcal{N} = 4$ SYM is conformal
- QCD demonstrates asymptotic freedom (coupling becomes weaker as energies increase to infinity), whereas $\mathcal{N} = 4$ SYM is strongly coupled at all length scales
- In QCD, both the fundamental and adjoint degrees of freedom are important to thermodynamic properties of QGP, whereas in $\mathcal{N} = 4$ SYM, there are no fundamental degrees of freedom

So, insights from hybrid model calculations in $\mathcal{N} = 4$ SYM are, and must be, treated **qualitatively and not quantitatively**.

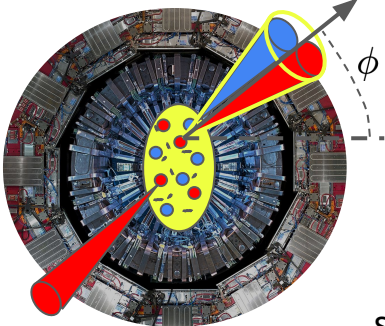
- $\mathcal{N} = 4$ $SU(N_c)$ SYM theory \longleftrightarrow IIB string theory in $AdS_5 \times S_5$
 - Cast particles as strings hanging from the 4-dimensional boundary into the 5-dimensional AdS bulk spacetime
 - Calculate observables of interest (ex: energy loss)
- Monte Carlo simulations of heavy ion collisions
 - Feed in some holographic calculations from above
 - Run the simulation and manipulate the output data to calculate more **observables that experimentalists can verify** using actual collider data



arXiv:1405.3864v3

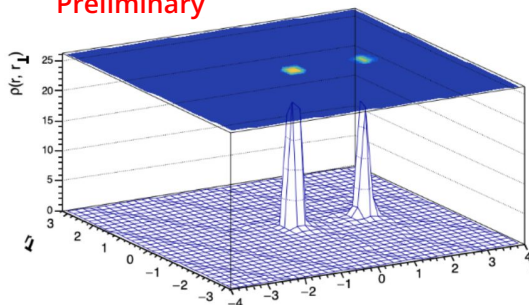
JET SUBSTRUCTURE

2-pronged substructure



Detector image from: CMS

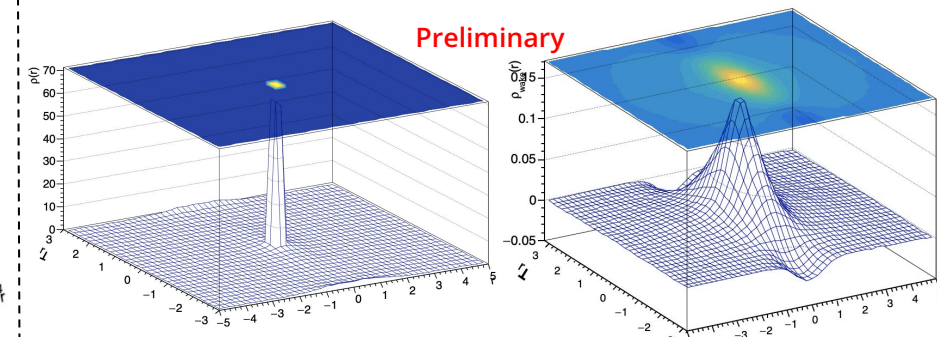
Preliminary



Shape of $R = 2.0$ anti-kt jet with 2 subjets separated by $1.8 < \Delta y < 2.0$ (Pb+Pb collisions)

JET-WAKES

Preliminary



Shape of $R = 0.2$ anti-kt jet (left) and its wake (right) in 0-5% centrality Pb+Pb collisions

TOPIC 3

Visualizing jet and wake substructure

Guiding question: When are the wakes left by different hard structures (jets) distinguishable?



Constructing Wide Jets from Skinny Subjets — ATLAS

[arXiv: 2301.05606] In Oct 2023, ATLAS published a paper that studied the substructure-dependence of large-radius jet suppression in Pb+Pb collisions at 5.02 TeV. They used the following procedure to reconstruct the large-radius jets:

- 1) Skinny jets with radius $R = 0.2$ were reconstructed using the anti- k_T algorithm
- 2) Large-radius $R = 1.0$ jets, restricted to $|y| < 2.0$, were reconstructed by clustering the $R = 0.2$ skinny (sub)-jets with $p_T > 35$ GeV and $|\eta| < 3.0$ using the anti- k_T algorithm with $R = 1.0$
- 3) The k_T algorithm was used to recluster the $R = 0.2$ subjets of each large-radius $R = 1.0$ jet, to define two substructure observables

$$\Delta R_{12} = \sqrt{\Delta y_{12}^2 + \Delta \phi_{12}^2} \quad \sqrt{d_{12}} = \min(p_{T1}, p_{T2}) \times \Delta R_{12}$$

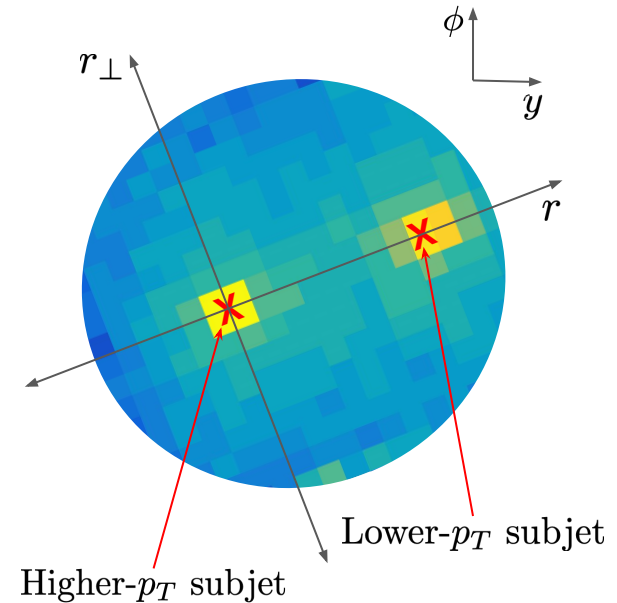
Respectively, these are the angular separation and k_T splitting scale between the two hardest constituents in the penultimate step of the k_T -reclustering

Example: If the large-radius jet had two $R = 0.2$ subjets, then these two subjets would be the two hardest constituents in question. If the large-radius jet had only 1 subjet, then

$$\Delta R_{12} = 0 \text{ and } \sqrt{d_{12}} = 0.$$

OUR COORDINATE SYSTEM

- We first constructed large-radius $R = 2.0$ jets from small-radius $R = 0.2$ subjets using the same procedure as ATLAS. We are using a larger radius of $R = 2.0$ to observe how the **substructure of jet-wakes differs between closely-separated and far-separated subjets**.
- Restrict to only $R = 2.0$ jets with 1 or 2 subjets
- Given a large-radius $R = 2.0$ jet with two subjets, we define a coordinate system (r, r_{\perp}) as such:
 - 1) Let the higher- p_T and lower- p_T subjets be located at $(y_{\text{high}}, \phi_{\text{high}})$ and $(y_{\text{low}}, \phi_{\text{low}})$, respectively.
 - 2) Define the origin of our new coordinates to be at $(y_{\text{high}}, \phi_{\text{high}})$ and define the r -axis to point positively in the direction of $(y_{\text{low}}, \phi_{\text{low}})$.
 - 3) Define the r_{\perp} -axis to be perpendicular to the r -axis, such that $\hat{r}_{\perp} = \hat{y} \times \hat{\phi}$.
- For an $R = 2.0$ jet with 1 subjet, the r -axis is centered on the y -coordinate of the subjet, and points in the y -direction. Similarly, the r_{\perp} -axis is centered on the ϕ -coordinate of the subjet, and points in the ϕ -direction



JET SHAPE OBSERVABLES

Now, we define the jet shape to measure the average fraction of a jet's transverse momentum within specified ranges of r and r_{\perp} . We can select large-radius jets in different ranges of ΔR_{12} (or Δy_{12}) to study the **dependence of jet shape on subjet-separation**.

For a specified range of ΔR_{12} (or Δy_{12}), we have the following jet shape observables:

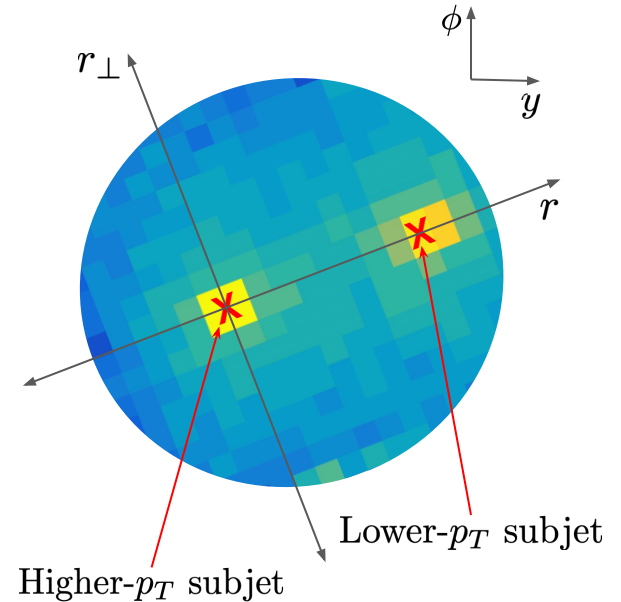
$$\rho^{2d}(r, r_{\perp}) = \frac{1}{\Delta R_{12}} \frac{1}{\Delta r} \frac{1}{\Delta r_{\perp}} \frac{1}{N_{\text{jet}}} \sum_{\text{jets}} \left(\frac{1}{p_T^{\text{jet}}} \left[p_T \right]_{(r - \frac{\Delta r}{2}, r_{\perp} - \frac{\Delta r_{\perp}}{2})}^{(r + \frac{\Delta r}{2}, r_{\perp} + \frac{\Delta r_{\perp}}{2})} \right)$$

$$\rho^{2d}(r, r_{\perp}) = \frac{1}{\Delta y_{12}} \frac{1}{\Delta r} \frac{1}{\Delta r_{\perp}} \frac{1}{N_{\text{jet}}} \sum_{\text{jets}} \left(\frac{1}{p_T^{\text{jet}}} \left[p_T \right]_{(r - \frac{\Delta r}{2}, r_{\perp} - \frac{\Delta r_{\perp}}{2})}^{(r + \frac{\Delta r}{2}, r_{\perp} + \frac{\Delta r_{\perp}}{2})} \right)$$

$$\rho^{1d}(r) = \text{Proj}_r(\rho^{2d}(r, r_{\perp}))$$

Important remarks:

- In our calculations of jet shapes, we include all particles within an $R = 2.0$ radius of the axis of each large-radius jet, not just the particles inside the $R = 0.2$ skinny subjets.
- When experimentalists measure jet shape, they have to subtract the background - we don't have to do this.
- **Only including hadrons from the wake in our calculation of jet shape allows us to plot the shape of the large-radius jet-wake.**



TOPIC 4

Results from gamma-jet calculations

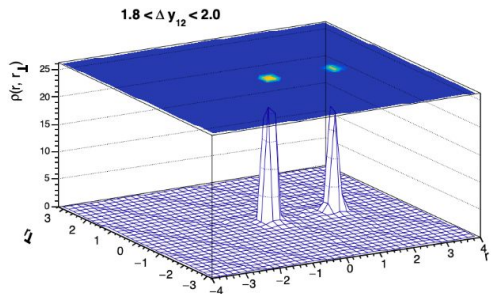
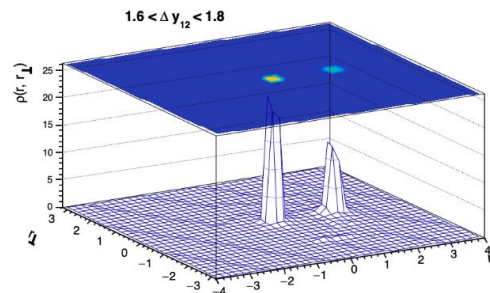
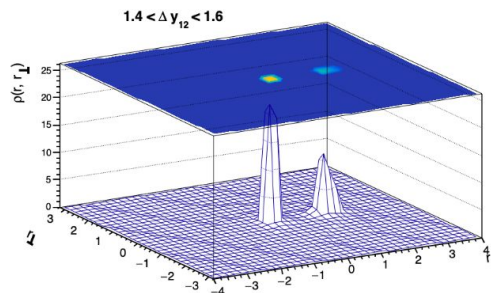
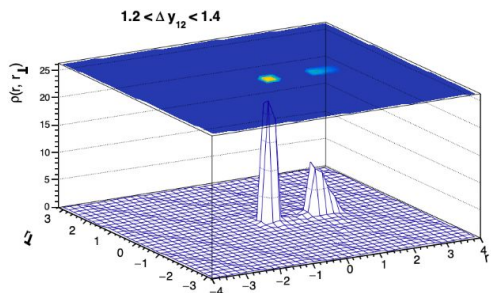
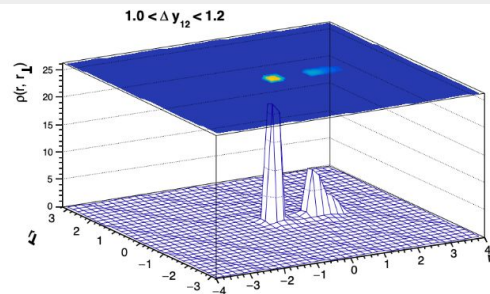
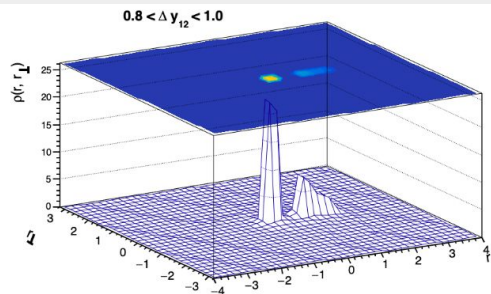
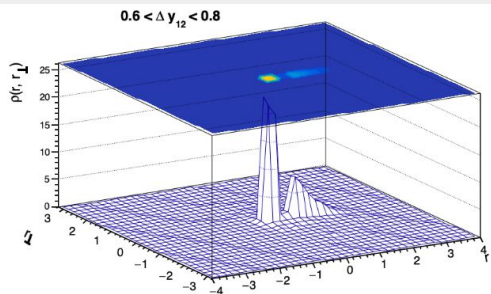
GAMMA-JETS IN Pb+Pb COLLISIONS

- We show results of the jet shapes calculated for gamma-jet events in Pb+Pb collisions. Gamma-jets are jet events where the recoiling jet is a photon.
- The **photon produces no wake of its own**, and so the jet shape will look cleaner than in the case of dijet events (events with almost back-to-back jets).
- Photons were selected using the following selection and isolation criteria:
 - $p_T^Y > 100$ GeV
 - $|\eta^Y| < 1.44$
 - The total transverse energy around a 0.4 radius of the photon must be less than 5 GeV
 - $\Delta\phi_{\gamma, \text{jet}} > 2\pi/3$

Important note: The photon is not considered a jet on its own in our analysis. So, **none of the photons contribute to the jet shape observables** we calculated.

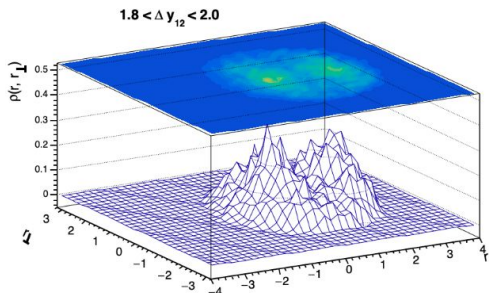
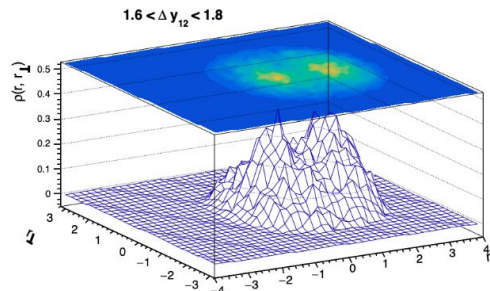
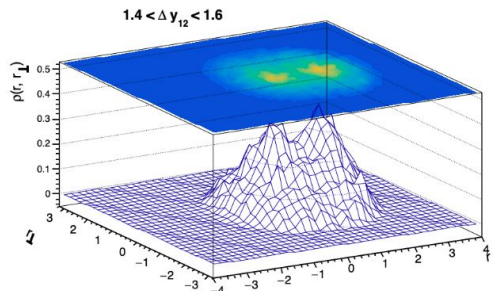
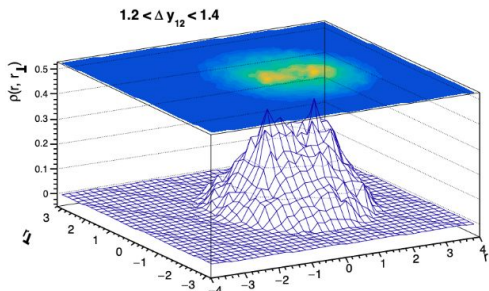
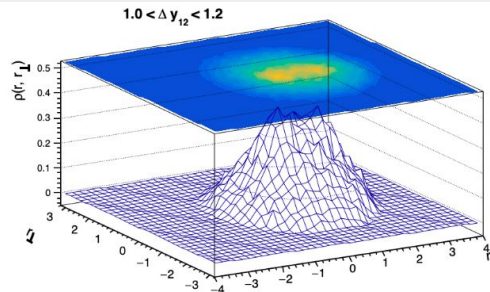
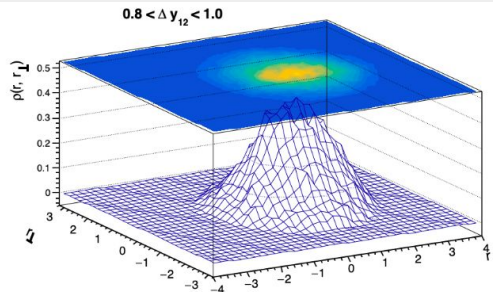
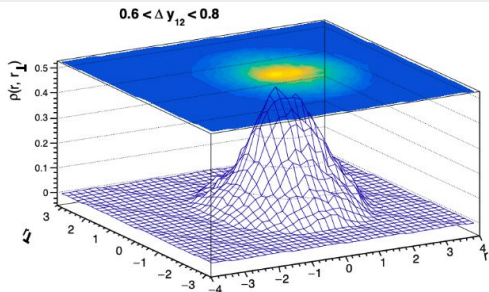
- Large radius $R = 2.0$ jets were constructed from $R = 0.2$ subjets whose $|p_T| > 35$ GeV and $|\eta| < 3.0$. In our analysis of gamma-jets, we restricted to examining only the large-radius jets with $|y| < 2.0$ and $50 < p_T < 1000$ GeV.

JET SHAPE OF GAMMA-JETS IN Pb+Pb COLLISIONS WITH 2 SUBJETS



They look as we expected...

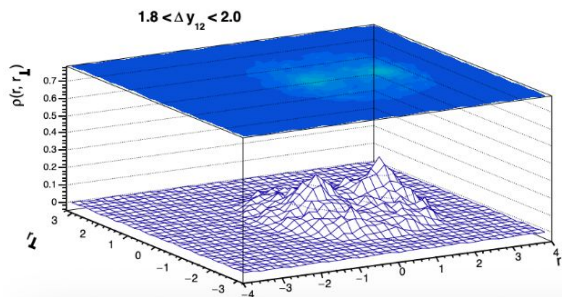
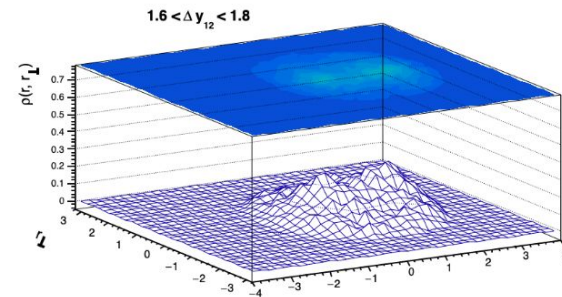
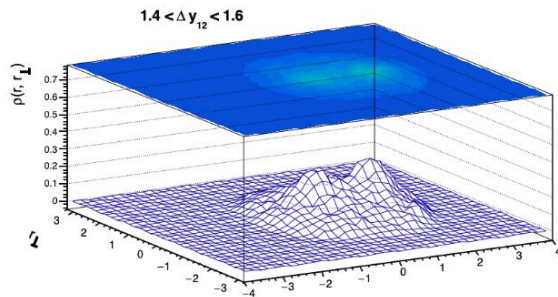
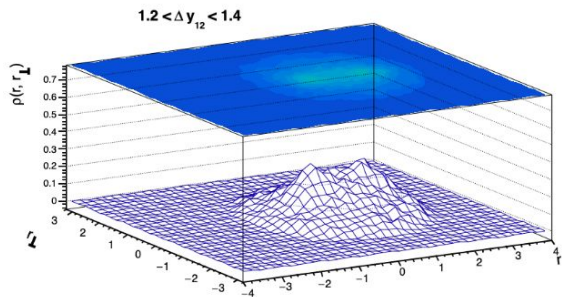
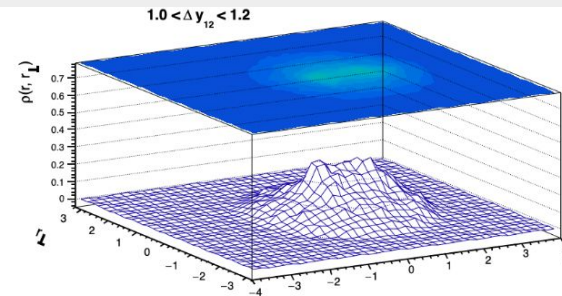
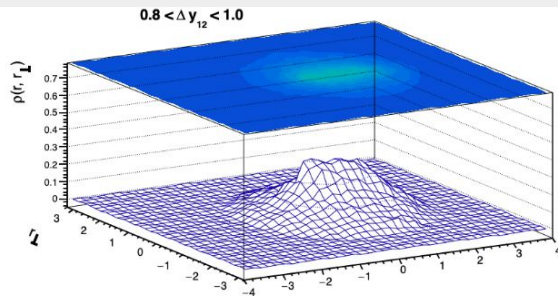
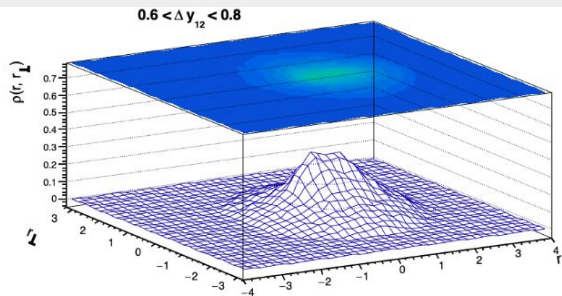
WAKE SHAPE OF GAMMA-JETS IN Pb+Pb COLLISIONS WITH 2 SUBJETS



For closely-separated subjects ($\Delta y_{12} < 1.0$), there is a single wake produced by 2 hard structures (the subjects). **Two distinct wakes are visibly produced only when the subjects are far-separated** (around $\Delta y_{12} > 1.2$)!

Can we see this in experiments?

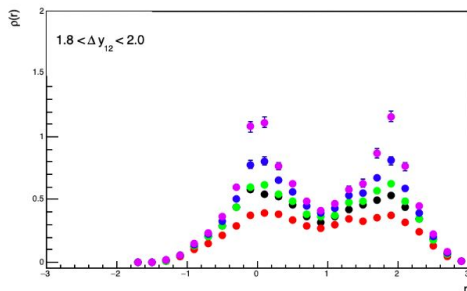
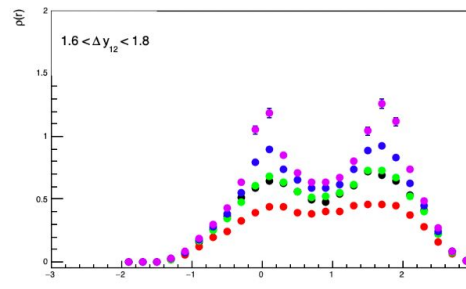
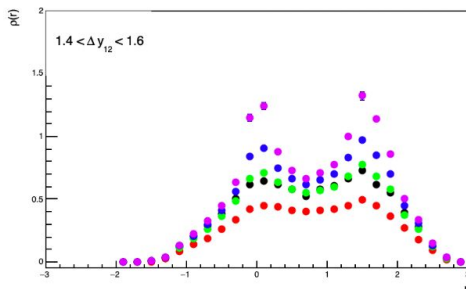
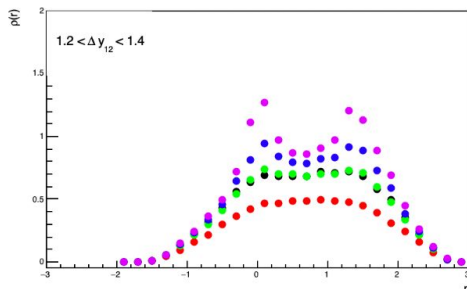
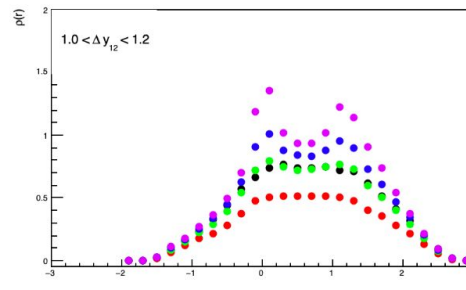
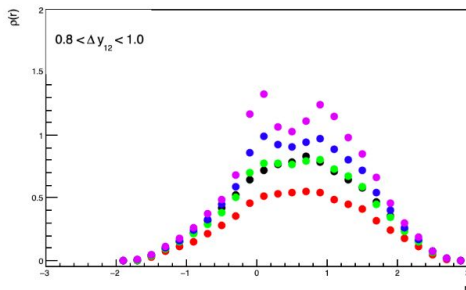
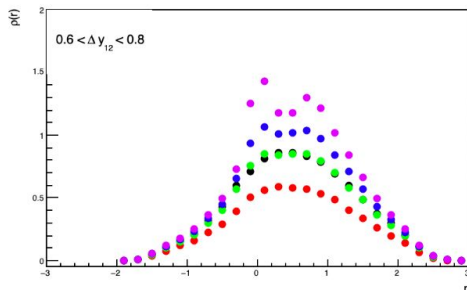
SHAPE OF PARTICLES WITH $p_T < 1.0$ GeV IN GAMMA-JETS IN Pb+Pb COLLISIONS WITH 2 SUBJETS



Looking at the shape of all particles (wake + nonwake) with $p_T < 1$ GeV acts as a good proxy of the 3-dimensional shape of the wake. The wake-substructures match!

But, experimentalists subtract the background when they calculate observables. This amounts to discounting particles with p_T below some cutoff (e.g. 0.7 GeV, [https://link.springer.com/article/10.1007/JHEP05\(2018\)006](https://link.springer.com/article/10.1007/JHEP05(2018)006)).

SHAPE OF WAKE AND SOFT PARTICLES GAMMA-JETS IN Pb+Pb COLLISIONS WITH 2 SUBJETS



- All particles with $p_T < 3.0$ GeV
- All particles with $p_T < 2.0$ GeV
- All particles with $p_T < 1.5$ GeV
- All particles with $p_T < 1.0$ GeV
- Only wake particles (any p_T)

The projection of the jet shape onto the r -axis for soft particles with $p_T < 1.5$ GeV most resembles the shape of the wake!

Perhaps restricting to particles with $0.7 < p_T < 1.5$ GeV will give us a good visual proxy for the wake in experiments!

SUMMARY

- Jets are useful phenomena to study the physics of QCD at high temperatures
- AdS/CFT is a useful tool to probe the strongly coupled, nonperturbative regimes of jet physics
- The substructure of the wake sourced by a jet's energy loss is nontrivially dependent on subjet-separation
 - When 2 subjets are closely-separated by $\Delta y_{12} < 1.0$ a single wake is produced, whereas when 2 subjets are far-separated (around $\Delta y_{12} > 1.4$), two distinct subwakes are produced.
- This nontrivial dependence of wake-substructure on jet-substructure illuminates a way to visualize jet-wakes in experiments

Remarks

- In this presentation, I have only included the results for our analysis of gamma-jets. The results for an analysis of inclusive jets yielded qualitatively similar conclusions.
- The results presented here are for the hybrid strong/weak coupling model, not for nature. **The main point is that our analysis presents methods and observables that experimentalists can use to see the wake and study wake-substructure!**



Website: <https://www.arjunkudinoor.com/talks/hep>

- These slides
- References to read (textbooks, papers)
- You can contact me

THANKS FOR LISTENING!

BACKUP SLIDES

dim-4 CFTs of gauge
(particle) fields



dim-5 gravitational theories in
AdS spacetime

Also called the **AdS/CFT correspondence**, first proposed by Juan Maldacena in 1997. A particular formulation conjectures an equivalence between the two descriptions of N_c D3-branes:

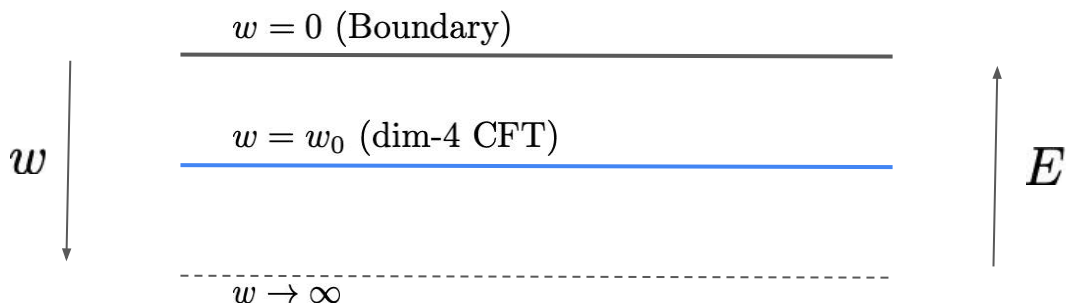
- 1) A hyperplane of flat spacetime with open strings attached. The low energy limit is described by an $\mathcal{N} = 4$ SYM theory with gauge group $SU(N_c)$.
- 2) A curved spacetime geometry where only closed strings propagate. The low energy limit is described by a closed Type IIB string theory in $AdS_5 \times S_5$.

Equating the low-energy limits gives a duality

$$\mathcal{N} = 4 \text{ } SU(N_c) \text{ SYM theory} \quad \longleftrightarrow \quad \text{IIB string theory in } AdS_5 \times S_5$$

Punchline: **Difficult calculations in strongly coupled gauge theories may be solved in its more tractable weakly coupled (small curvature) gravitational dual**; vice versa. The planar (large N_c) strongly coupled (large $\lambda = g_{\text{YM}}^2 N_c$) limit of the SYM theory can be described using just a classical supergravity dual.

Suppose we have a 4-dimensional non-gravitational Minkowski CFT at energy scale $E_0 \sim \frac{1}{w_0}$, where w is a 5th space-like dimension. So, we can describe this 4-dimensional gauge theory at all energy scales as a 5-dimensional **holographic** theory with coordinates (t, x, y, z, w) .



The most general (4+1)-dimensional metric that is consistent with Poincaré in 4 dimensions is

$$ds^2 = \Omega^2(w)(-dt^2 + dx^2 + dy^2 + dz^2 + dw^2).$$

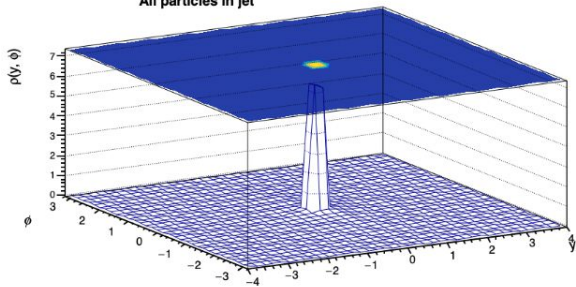
A CFT is invariant under rescalings $(t, x, y, z) \mapsto C(t, x, y, z)$. So, $\Omega(w) \mapsto \frac{1}{C}\Omega(w)$ under $w \mapsto Cw$.

Thus, $\Omega(w) = \frac{R}{w}$ and $ds^2 = \frac{R^2}{w^2}(-dt^2 + dx^2 + dy^2 + dz^2 + dw^2)$.

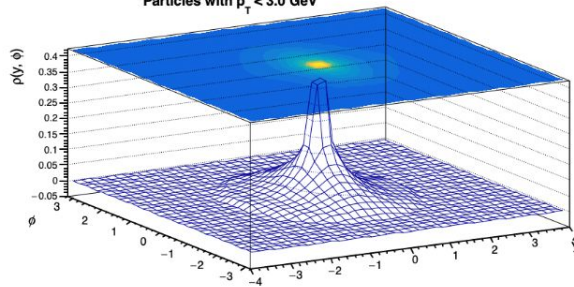
This is the metric of **(4+1)-dimensional AdS spacetime** with constant negative curvature proportional to $\frac{1}{R^2}$, i.e. this 5-dimensional spacetime has **gravity!**

SHAPES OF GAMMA-JETS IN Pb+Pb COLLISIONS WITH 1 SUBJET

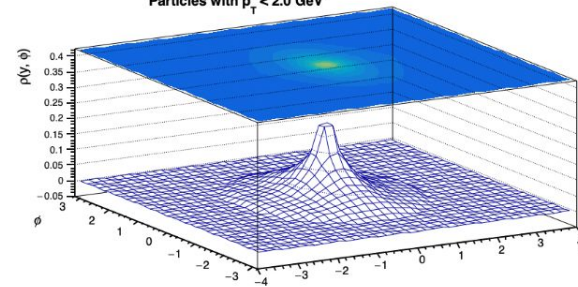
All particles in jet



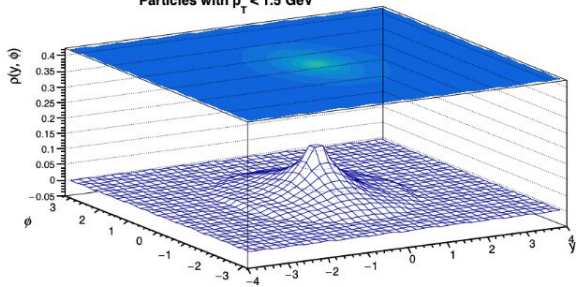
Particles with $p_T < 3.0$ GeV



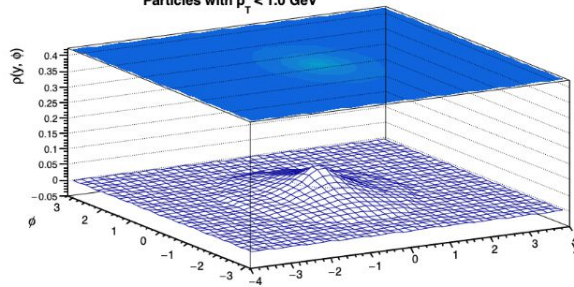
Particles with $p_T < 2.0$ GeV



Particles with $p_T < 1.5$ GeV



Particles with $p_T < 1.0$ GeV



Wake particles only (any p_T)

